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**AN ALGORITHM FOR THE BASIS
OF THE FINITE FOURIER TRANSFORM**

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Abstract

The Finite Fourier Transformation matrix (F.F.T.) plays a central role in the formulation of quantum mechanics in a finite dimensional space studied by the author over the past couple of decades. An outstanding problem which still remains open is to find a complete basis for F.F.T. In this paper we suggest a simple algorithm to find the eigenvectors of F.F.T.

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I. INTRODUCTION

The finite Fourier transform matrix (F.F.T.) plays a fundamental role in many contexts and has been studied extensively [1-3]. It is central in the discussions on finite dimensional quantum mechanics based on Weyl's commutation relations [4] studied by the author in a series of publications [5]. The eigenvalues of this matrix were determined by Schur [1] and a simple argument to recover this result has been given earlier [6]. The calculation of the eigenvectors is not straightforward and many methods have been given in particular, by Mehta [7]. In Section IV, we present a new algorithm to find the eigenvectors.

II. EIGENVALUES OF S

The F.F.T. matrix S , which is unitary, is defined by

$$S_{\alpha\beta} = \frac{1}{\sqrt{n}} \exp \left[\frac{2\pi i}{n} \alpha \beta \right],$$

$$\alpha, \beta = 0, 1, 2, \dots, n-1 \quad (2.1)$$

$$i = \sqrt{-1}$$

and has many interesting properties

$$1) \quad (S^2)_{\alpha\beta} \equiv I'_{\alpha\beta} = \delta_{\alpha + \beta, 0} \pmod{n} \quad (2.2)$$

Since $S^2 f_{\alpha} = f_{-\alpha \pmod{n}}$, for a vector f_{α} with n components, S^2 is called the parity operator

$$2) \quad (S^4)_{\alpha\beta} = \delta_{\alpha\beta} \quad (2.3)$$

like the usual Fourier transform.

3) The matrix S , which is by definition a symmetric matrix will diagonalize any circulant matrix.

From Equation (2-3), it is clear that the eigenvalues of S are simply ± 1 and $\pm i$. There is then a degeneracy of the eigenvalues. The first problem will be to determine this. Luckily, Equations (2.1)-(2.3) can be repeatedly used to fix this [6]. If k_1, k_2, k_3 and k_4 denote the multiplicity of the eigenvalues taken in the order (1, -1, i, -i), Equation (2.1) implies that

$$\begin{aligned} \text{Tr } S &= \frac{1}{\sqrt{n}} \sum_{\ell=0}^{n-1} \left[\exp \frac{2\pi i \ell}{n} \right] \ell^2 \\ &= \frac{1}{2} (1 + i) \left[1 + \exp \left(\frac{-i\pi n}{2} \right) \right], \end{aligned} \tag{2.4}$$

and hence

$$\begin{aligned} \text{Tr } S &= (k_1 - k_2) + i(k_3 - k_4) \\ &= 1 \text{ for } n = 4k + 1, \\ &= 0 \text{ for } n = 4k + 2, \\ &= i \text{ for } n = 4k + 3, \\ &= (1 + i) \text{ for } n = 4k, \\ k &= 0, 1, 2, \dots \end{aligned} \tag{2.5}$$

From Equation (2) we infer that

$$\begin{aligned} \text{Tr } S^2 &= (k_1 + k_2) - (k_3 + k_4) \\ &= 1 \text{ for } n \text{ odd,} \\ &= 2 \text{ for } n \text{ even.} \end{aligned}$$

We also have

$$\text{Tr } S^4 = n = k_1 + k_2 + k_3 + k_4.$$

Equations (2.5), (2.6) and (2.7) can be used to solve for k_1, k_2, k_3 and k_4 and one finds that

	$n = 4k + 1$	$n = 4k + 2$	$n = 4k + 3$	$n = 4k$
k_1	$k + 1$	$k + 1$	$k + 1$	$k + 1$
k_2	k	$k + 1$	$k + 1$	k
k_3	k	k	$k + 1$	k
k_4	k	k	k	$k - 1$

III. EIGENVECTORS OF S

Let us decompose S into its primitive idempotents as

$$S = \sum_{j=1}^4 i^j B(j),$$

where

$$\begin{aligned} B(1) &= \frac{1}{2}s + \frac{1}{4}(I - I'), \\ B(2) &= -\frac{1}{2}c + \frac{1}{4}(I + I'), \\ B(3) &= -\frac{1}{2}s + \frac{1}{4}(I - I'), \\ B(4) &= \frac{1}{2}c + \frac{1}{4}(I + I'), \end{aligned} \tag{3.2}$$

$$C_{\alpha\beta} = \frac{1}{\sqrt{n}} \cos \left(\frac{2\pi}{n} \alpha\beta \right)$$

$$s_{\alpha\beta} = \frac{1}{\sqrt{n}} \sin \left(\frac{2\pi}{n} \alpha\beta \right),$$

$$\alpha, \beta = 0, 1, 2, \dots, n-1 \tag{3.3}$$

It is easily verified that

$$S B(j) = i^j B(j), \tag{3.4}$$

thus the nonzero columns of B(j) yield the eigenvectors of S with eigenvalue i^j . Also, in analogy with the standard case, Mehta [7] has been able to express these eigenvectors in terms of Hermite functions with

discrete arguments.

IV. EIGENVECTORS OF S; AN ALTERNATE METHOD

Since the F.F.T. matrix S satisfies Equation (2.1) we construct the matrix [10]

$$\begin{aligned} T &= S^3 + S^2 S_d + S S_d^2 + S_d^3 \\ &= I' (S + S_d) + (S + S_d) S_d^2, \end{aligned} \quad (4.1)$$

where

$$S_d = \text{diagonal } S. \quad (4.2)$$

We find that

$$\begin{aligned} S T &= S (S^3 + S^2 S_d + S S_d^2 + S_d^3) \\ &= (I + S^3 S_d + S^2 S_d^2 + S S_d^3) \\ &= (S_d^3 + S^3 + S^2 S_d + S S_d^2) S_d \\ &= T S_d. \end{aligned} \quad (4.3)$$

If T is nonsingular,

$$T^+ S T = S_d \quad (4.4)$$

Therefore, the columns of T automatically provide the eigenvectors of S. The degenerate eigenvectors of S corresponding to the repeated eigenvalues can be made orthonormal by using Gram-Schmidt process. This will render T unitary. While the process is quite general, we shall illustrate this for some special cases

case of $n = 2$

$$S = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}, \quad (4.5)$$

and

$$S_d = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \quad (4.6)$$

$$\text{Since } S^2 = S_d^2 = I, \quad (4.7)$$

We get from Equation (4.1)

$$\begin{aligned} T &= 2 (S + S_d), \\ &= 2 \begin{pmatrix} 1 + \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & -1 - \frac{1}{\sqrt{2}} \end{pmatrix} \end{aligned} \quad (4.8)$$

We unitarized matrix of the eigenvectors of S is therefore

$$U_2 = \frac{1}{\sqrt{2\sqrt{2}(\sqrt{2}+1)}} \begin{pmatrix} \sqrt{2} + 1 & 1 \\ 1 & -(\sqrt{2} + 1) \end{pmatrix} \quad (4.9)$$

case of $n = 3$

$$\begin{aligned} S &= \frac{1}{\sqrt{3}} \begin{pmatrix} 1 & 1 & 1 \\ 1 & \epsilon & \epsilon^2 \\ 1 & \epsilon^2 & \epsilon \end{pmatrix}, \\ \epsilon &= \exp \frac{2\pi i}{3}. \end{aligned} \quad (4.10)$$

From Equation (2.8) we see that

$$S_d = \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & i \end{pmatrix} \quad (4.11)$$

one finds from Equation (4.1) that the unitarized matrix of the eigenvectors of S is

$$U_3 = \frac{1}{\sqrt{2} \sqrt{3} (\sqrt{3} + 1)} \begin{pmatrix} \sqrt{3} + 1 & \sqrt{2} & 0 \\ 1 & \frac{1 + \sqrt{3}}{\sqrt{2}} & i\sqrt{3 + \sqrt{3}} \\ 1 & \frac{1 + \sqrt{3}}{\sqrt{2}} & -i\sqrt{3 + \sqrt{3}} \end{pmatrix}. \quad (4.12)$$

case of n = 4

In this case we have

$$S = \frac{1}{2} \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & i & -1 & -i \\ 1 & -1 & 1 & -i \\ 1 & -i & -1 & i \end{pmatrix} \quad (4.13)$$

and

$$S_d = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & i \end{pmatrix} \quad (4.14)$$

It is easily calculated that

$$T = \begin{pmatrix} 3 & 1 & 1 & 0 \\ 1 & 1 & -1 & 2i \\ 1 & -1 & -1 & 0 \\ 1 & 1 & -1 & -2i \end{pmatrix} \quad (4.15)$$

The first two column vectors correspond to the eigenvalue = +1, the third one to -1 and the last to -i.

By a simple use of Gram-Schmidt orthogonalization procedure one can find the unitarized matrix corresponding to the eigenvectors of S as

$$U_4 = \frac{1}{\sqrt{2}\sqrt{4}(\sqrt{4}+1)} \begin{pmatrix} 3 & 0 & \sqrt{3} & 0 \\ 1 & \sqrt{2} & -\sqrt{3} & i\sqrt{6} \\ 1 & -2\sqrt{2} & -\sqrt{3} & 0 \\ 1 & \sqrt{2} & -\sqrt{3} & -i\sqrt{6} \end{pmatrix} \quad (4.16)$$

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REFERENCES

- [1] I. Schur, Nachr K. Gesell, Wiss. Gobbigen, Math. Phys. K1, 1921, 147; E. Landau, Elementary Number Theory (Chelsea, New York, 1966), p. 207-212.
- [2] B. C. Berndt and R. J. Evans, Bull. Am. Math. Soc. 5, 1981, 107; see also, Bull. Am. Math. Soc. 7, 1982, 441.
- [3] I. J. Good, Am. Math. Month. 69, 1962, 259; L. Anslander and R. Tolimieri, Bull. Am. Math. Soc. 1, 1979, 847; P. Morton, J. Num. Th., 12, 1980, 122; R. Tolimieri, Adv. Appl. Math. 5, 1984, 56; Allad. Ramakrishnon, 1972, L. Matrix Theory and the Grammar of Dirac Matrices (Tata McGraw Hill Publishing Co.).
- [4] H. Weyl, Theory of Groups and Quantum Mechanics, (Dover, New York) 1931, p. 272-280.
- [5] T. S. Santhanam and A. R. Tekumalla, Found. Phys. 6, 1976; T. S. Santhanam, Uncertainty Principle and Foundations of Qunatum Mechanics, Eds. W. C. Price and S. S. Chissick; (John Wiley and Sons) 1982, p. 227-243, Physica 114A, 445; R. Jagannalhan, T. S. Santhanam and R. Vasudevan, Int. J. Theor. Phys. 20, 1981, p. 755.
- [6] T. S. Santhanam and S. Madivanane, preprint IC/82/12, ICTP, Trieste, Italy (unpublished).
- [7] M. L. Mehta, J. Math Phys. 7, 1987, p. 781.
- [8] J. M. Jauch, Foundations of Quantum Mechanics (Addison-Wesley, Reading, MA), Sec. 13.7. See also, K. B. Wolf, 1972, Integral Transforms in Science and Engineering, (Plenum Press, NY).
- [9] See, for instance, F. R. Gantmacher, Matrix Theory (Chelsea, New York, 1959), Vol. I, p. 239.
- [10] This can be easily generalized to $A^{n-1} + A^{n-2} A_d + A^{n-3} A_d^2 + \dots + A A_d^{n-2} + A_d^{n-1}$ for the case of a general involution matrix satisfying the relation $A^n = I$.